

# An unfitted $hp$ -interface penalty finite element method for elliptic interface problems

Haijun Wu\* Yuanming Xiao†

## Abstract

An  $hp$  version of interface penalty finite element method ( $hp$ -IPFEM) is proposed for elliptic interface problems in two and three dimensions on unfitted meshes. Error estimates in broken  $H^1$  norm, which are optimal with respect to  $h$  and suboptimal with respect to  $p$  by half an order of  $p$ , are derived. Both symmetric and non-symmetric IPFEM are considered. Error estimates in  $L^2$  norm are proved by the duality argument.

**Key words.** Elliptic interface problems, unfitted mesh,  $hp$ -IPFEM

**AMS subject classifications.** 65N12, 65N15, 65N30

## 1 Introduction

Let  $\Omega = \Omega_1 \cup \Gamma \cup \Omega_2$  be a bounded and convex polygonal or polyhedral domain in  $\mathbb{R}^d$ ,  $d = 2$  or  $3$ , where  $\Omega_1$  and  $\Omega_2$  are two subdomains of  $\Omega$  and  $\Gamma = \partial\Omega_1 \cap \partial\Omega_2$  is a  $C^2$ -smooth interface (see Fig. 1). Consider the following elliptic interface problem:

$$(1) \quad \begin{cases} -\nabla \cdot (a(x)\nabla u) = f, & \text{in } \Omega_1 \cup \Omega_2, \\ [u] = g_D, \quad [(a(x)\nabla u) \cdot \mathbf{n}] = g_N, & \text{on } \Gamma, \\ u = 0, & \text{on } \partial\Omega, \end{cases}$$

where  $[v] = v|_{\Omega_1} - v|_{\Omega_2}$  denotes the jump of  $v$  across the interface  $\Gamma$ ,  $\mathbf{n}$  is the unit outward normal to the boundary of  $\Omega_1$ , and  $a(x)$  is bounded from below and above by some positive constants. Note that  $a(x)$  is allowed to be discontinuous across the interface  $\Gamma$ .

It is well known that the finite difference method (FDM) compared to the finite element method (FEM) is easy to implement since it discretizes the PDEs on simple meshes. But it is not easy to construct high order FDMs when the

---

\*Department of Mathematics, Nanjing University, Jiangsu, 210093, P.R. China. ([hjw@nju.edu.cn](mailto:hjw@nju.edu.cn)). The work of this author was partially supported by the national basic research program of China under grant 2005CB321701 and by the NSF of China grant 10971096.

†Department of Mathematics, Nanjing University, Jiangsu, 210093, P.R. China. ([xym@nju.edu.cn](mailto:xym@nju.edu.cn)).

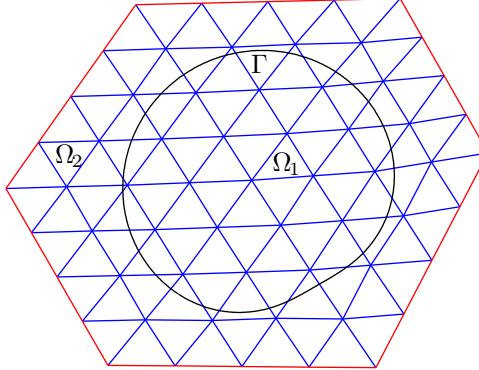


Figure 1: A sample domain  $\Omega$  and an unfitted mesh.

geometries are complicated. By contrast, the FEM provides a systematic way to design high order schemes for problems with complicated geometries, while the classical FEMs require that the interfaces and boundaries are fitted well enough by underlying meshes in order to achieve optimal convergence rates (cf. [7, 12, 27]). Therefore it has always been attractive to develop systematic ways for constructing high order numerical methods on simple unfitted meshes. Unfitted mesh techniques are particularly convenient for time-dependent problems when moving interfaces are involved. One could just use the same mesh on the domain for different times instead of repeatedly remeshing the domain to fit the moving interfaces. Another motivation for developing unfitted methods is from the adaptive finite element methods. We know that in an adaptive finite element procedure, the meshes are locally refined repeatedly to equidistribute the error. But it is readily to encounter the “inverted elements”, i.e., the elements with negative directional areas/volumes (cf. [37]), near curved interfaces during interface-fitted local refinements, in particular, in the three dimensional case.

Various finite difference schemes have been proposed on unfitted Cartesian grids, for example, the *immersed boundary method* by Peskin [34], the *ghost fluid method* by Osher and his coworkers [17], the *immersed interface method* by LeVeque and Li [26], and the *matched interface and boundary method* by Zhou, et al. [41]. These methods have the first order, second order or higher order truncation errors near the interfaces. They have been generalized in various ways and applied to many problems [30, 31, 38, 32, 24, 25, 29, 40, etc].

The existing finite element schemes on unfitted meshes are usually designed

by proper modifications of the standard finite element methods near the interfaces. The *immersed finite element method* by Li, et al. [28, 19] modifies the basis functions for nodal points near the interfaces to satisfy the homogeneous jump conditions. The linear immersed finite element method is optimally convergent in  $H^1$  and  $L^2$  norms for two dimensional problems [13]. We refer to [11] for *adaptive immersed interface finite element methods* for solving elliptic and Maxwell interface problems with singularities. The *multiscale finite element method* by Chu, Graham, and Hou [14] modifies the basis functions near the interface by solving appropriately designed subgrid problems. The method was proved, for the two dimensional case, to be of first-order accuracy in  $H^1$  norm and of second-order accuracy in  $L^2$  norm. Note that when the interface intersects an element in a straight line, the multiscale finite element method coincides with the linear immersed finite element method (cf. [14]). The *penalty finite element method* proposed by Babuška [4] modifies the bilinear form near the interface by penalizing the jump of the solution value across the interface. The method may use high order elements but is suboptimally convergent in  $h$ . The linear case of the method was analyzed again by Barrett and Elliott [7] and was proved to be optimally convergent in  $h$  under more regularity assumptions on the exact solution. We remark that the idea of adding penalty terms is widely used in the interior penalty Galerkin methods [2, 3, 6, 16, 18, 35, etc]. The *unfitted finite element method* proposed by A. Hansbo and P. Hansbo [23] can be viewed as an improvement of the linear version of the Babuška's method. This unfitted method was proved, for two dimensional elliptic interface problems, to be optimally  $h$ -convergent in both  $H^1$  and  $L^2$  norms and the error estimates are uniform with respect to the relative position of the interface to the unfitted grid. The key idea of their method is using weighted average flux across the interface in the bilinear form instead of using the arithmetic average one as in the standard interior penalty discontinuous Galerkin methods. In their method, the weights along the intersection of the interface and a grid element are chosen according to the areas of the two parts of the element separated by the interface. More recently, Massjung [33] proposed an *hp-unfitted discontinuous Galerkin method* for Problem (1) which uses weighted average flux across the interface and penalizes not only the jump condition on the solution value but also the jump condition on the flux. It was proved that, for the two dimensional case, the method converges in broken  $H^1$  norm at an optimal rate with respect to  $h$  and at a suboptimal rate with respect to  $p$  by a factor of  $p$ . The flux penalty term is not essential in the convergence analysis but helpful for numerical stability (cf. [33]). We also refer to [8, 22] for recent more studies on the discontinuous Galerkin methods for elliptic interface problems.

In this paper we propose an *hp*-interface penalty finite element method (*hp*-IPFEM) for the interface problem (1). In this method we penalize the jump conditions on the solution value as well as the flux across the interface and adopt the arithmetic average flux across the interface in the bilinear form. We provide a rigorous analysis to show that, for both two and three dimensional cases, the error estimates in broken  $H^1$  norm are optimal with respect to  $h$  and suboptimal with respect to  $p$  by a half order of  $p$ . Both symmetric and non-symmetric

IPFEM are considered. Some  $hp$ -error estimates in  $L^2$  norm are also derived by the duality argument. We would like to mention that our tricks on dealing with discontinuities across the interface are common in the standard interior penalty discontinuous Galerkin methods. They are adopted by [23, 33] except the arithmetic average one. Sure we do not add penalty terms on edges/faces of elements away from the interface as the standard interior penalty discontinuous Galerkin methods do.

For any intersection segment/patch  $e$  between the interface and a grid element  $K^e$ , let  $K_i^e = K^e \cap \Omega_i, i = 1, 2$ , be the two parts of  $K^e$  separated by the interface. The key idea of our analysis is the use of the following straight forward identity on  $e$ :

$$\frac{v_1 + v_2}{2} = v_i + \frac{(-1)^i}{2}(v_1 - v_2), \quad \text{where } v_i = (a\nabla v_h)|_{K_i^e} \cdot \mathbf{n}, \quad i = 1, 2.$$

By this we may bound the term of the arithmetic average flux (i.e.  $\frac{v_1 + v_2}{2}$ ) by any component in the average (say  $v_i$ ,  $i$  = either 1 or 2) and the jump of flux (i.e.  $v_1 - v_2$ ) which is controlled by the flux penalty term. This component is then estimated by a local inverse trace inequality on the corresponding part of the interface element. As a result, it suffices to prove the local inverse trace inequality on either  $K_1^e$  or  $K_2^e$  instead of on both of them as treated in [33]. Therefore we avoid discussing various relative positions of the interface to an element  $K^e$  which are already complicated in two dimensional case since  $K_1^e$  or  $K_2^e$  may not be shape regular. This makes our analysis much simpler and easy to apply to the three dimensional case. To the best of our knowledge, our results give the first  $hp$  a priori error estimates for interface problems in three dimensions and the best in two dimensions.

The rest of our paper is organized as follows. We formulate the  $hp$ -interface penalty finite element methods in Section 2 and list some preliminary lemmas in Section 3. The  $H^1$ - and  $L^2$ - error estimates of the symmetric interface penalty finite element methods are given in Section 4. The error estimates of the non-symmetric interface penalty finite element methods are given in Section 5. Section 6 is devoted to the proofs of the local trace inequality and the local inverse trace inequality.

Throughout the paper,  $C$  is used to denote a generic positive constant which is independent of  $h$ ,  $p$ , and the penalty parameters. We also use the shorthand notation  $A \lesssim B$  and  $B \gtrsim A$  for the inequality  $A \leq CB$  and  $B \geq CA$ .  $A \approx B$  is for the statement  $A \lesssim B$  and  $B \lesssim A$ . Note that the constants in this paper may depend on the jump of the coefficient  $a(x)$  across the interface. Although it is an interesting topic to derive error estimates with explicit dependence on the jump of the coefficient  $a(x)$ , it may bury our basic idea in a more complicated analysis. We leave this issue to future work.

## 2 Formulation of *hp*-interface penalty finite element methods

To formulate our *hp*-IPFEMs, we need to introduce some notation. The space, norm and inner product notation used in this paper are all standard, we refer to [9, 15] for their precise definitions.

Let  $\{\mathcal{M}_h\}$  be a family of conforming, quasi-uniform, and regular partitions of the domain  $\Omega$  (cf. [9]) into triangles and parallelograms/tetrahedrons and parallelepipeds. For any  $K \in \mathcal{M}_h$ , we define  $h_K := \text{diam}(K)$ . Let  $h := \max_{K \in \mathcal{M}_h} h_K$ . Then  $h_K \asymp h$ . Note that any element  $K \in \mathcal{M}_h$  is considered as closed. Obviously, each partition  $\mathcal{M}_h$  induces a partition, denoted by  $\mathcal{I}_h$ , of the interface  $\Gamma$  (cf. Fig. 1). Note that any interface segment/patch  $e \in \mathcal{I}_h$  is either contained entirely in one element in  $\mathcal{M}_h$  and has nonempty intersection with its interior, or is the common edge/face of two neighboring elements in  $\mathcal{M}_h$ . For any  $e \in \mathcal{I}_h$ , let  $K^e \in \mathcal{M}_h$  be one of the element(s) containing  $e$  and let  $K_i^e = K^e \cap \Omega_i, i = 1, 2$ .

Introduce the “energy” space

$$(2) \quad V := \{v : v|_{\Omega_i} = v_i, \text{ where } v_i \in H_0^1(\Omega), v_i|_K \in H^2(K), \forall K \in \mathcal{M}_h, i = 1, 2\}.$$

For any  $K \in \mathcal{M}_h$ , let  $\mathcal{P}_p(K)$  denote the set of all polynomials whose degrees in all variables (total degrees) do not exceed  $p$  if  $K$  is a triangle/tetrahedron, and the set of all polynomials whose degrees in each variable (separate degrees)  $\leq p$  if  $K$  is a parallelogram/parallelepiped. Denote by  $U_h^p$  the *hp*-continuous finite element space, that is,

$$(3) \quad U_h^p := \{v_h \in H_0^1(\Omega) : v_h|_K \in \mathcal{P}_p(K), \forall K \in \mathcal{M}_h\}.$$

We define our interface penalty finite element approximation space  $V_h^p$  as

$$(4) \quad V_h^p := \{v_h : v_h|_{\Omega_i} = v_{ih}, \text{ where } v_{ih} \in U_h^p, i = 1, 2\}.$$

Clearly,  $V_h^p \subset V \subset L^2(\Omega)$ . But  $V_h^p \not\subset H^1(\Omega)$ .

We also define the jump  $[v]$  and average  $\{v\}$  of  $v$  on the interface  $\Gamma$  as

$$(5) \quad [v] := v|_{\Omega_1} - v|_{\Omega_2}, \quad \{v\} := \frac{v|_{\Omega_1} + v|_{\Omega_2}}{2}.$$

Testing the elliptic problem (1) by any  $v \in V$ , using integration by parts, and using the identity  $[vw] = \{v\}\{w\} + [v]\{w\}$ , we obtain

$$\sum_{i=1}^2 \int_{\Omega_i} a \nabla u \cdot \nabla v - \int_{\Gamma} \{a \nabla u \cdot \mathbf{n}\} [v] - \int_{\Gamma} g_N \{v\} = \int_{\Omega} fv$$

Define the bilinear form  $a_h(\cdot, \cdot)$  on  $V \times V$ :

$$(6) \quad a_h(u, v) := \sum_{i=1}^2 \int_{\Omega_i} a \nabla u \cdot \nabla v - \sum_{e \in \mathcal{I}_h} \int_e \left( \{a \nabla u \cdot \mathbf{n}\} [v] + \beta [u] \{a \nabla v \cdot \mathbf{n}\} \right) + J_0(u, v) + J_1(u, v),$$

$$(7) \quad J_0(u, v) := \sum_{e \in \mathcal{I}_h} \frac{\gamma_0 p^2}{h_{K^e}} \int_e [u] [v],$$

$$(8) \quad J_1(u, v) := \sum_{e \in \mathcal{I}_h} \frac{\gamma_1 h_{K^e}}{p^2} \int_e [a \nabla u \cdot \mathbf{n}] [a \nabla v \cdot \mathbf{n}],$$

where  $\beta$  is a real number,  $\gamma_0$  and  $\gamma_1$  are nonnegative numbers to be specified later. Define further the linear form  $F_h(\cdot)$  on  $V$ :

$$(9) \quad F_h(v) := \int_{\Omega} f v + \int_{\Gamma} g_N \{v\} - \beta \int_{\Gamma} g_D \{a \nabla v \cdot \mathbf{n}\} + J_D(v) + J_N(v),$$

$$(10) \quad J_D(v) := \sum_{e \in \mathcal{I}_h} \frac{\gamma_0 p^2}{h_{K^e}} \int_e g_D [v],$$

$$(11) \quad J_N(v) := \sum_{e \in \mathcal{I}_h} \frac{\gamma_1 h_{K^e}}{p^2} \int_e g_N [a \nabla v \cdot \mathbf{n}].$$

It is easy to check that the solution  $u$  to the problem (1) satisfies the following formulation:

$$(12) \quad a_h(u, v) = F_h(v), \quad \forall v \in V.$$

We are now ready to define our  $hp$ -interface penalty finite element methods inspired by the formulation (12): Find  $u_h \in V_h^p$  such that

$$(13) \quad a_h(u_h, v_h) = F_h(v_h), \quad \forall v_h \in V_h^p.$$

**Remark 2.1.** (i) *The above tricks on dealing with the discontinuities are from the interior penalty discontinuous or continuous Galerkin methods (see, e.g., [2, 18, 10]). When  $\beta = 1$ ,  $a_h(\cdot, \cdot)$  is symmetric, that is,  $a_h(v, w) = a_h(w, v)$ , the method is termed as the symmetric interface penalty finite element method (SIPFEM), which corresponds to the symmetric interior penalty Galerkin method [2, 39]. On the other hand, when  $\beta \neq 1$ ,  $a_h(\cdot, \cdot)$  is non-symmetric. In particular, when  $\beta = -1$ , the method is called the non-symmetric interface penalty finite element method (NIPFEM), which would correspond to the non-symmetric interior penalty Galerkin method studied in [35]. In this paper, for the ease of presentation, we only consider the case  $\beta = 1$  and  $\beta = -1$ , nevertheless the ideas of this paper also apply to the case  $\beta \neq \pm 1$ .*

(ii) *The unfitted finite element method given in [23] does not include the penalty term  $J_1$  and uses a weighted average  $\{v\} = \kappa_1 v|_{\Omega_1} + \kappa_2 v|_{\Omega_2}$  instead of the arithmetic one in (5), where  $\kappa_i|_e = |K_i^e| / |K|$ . Optimal error*

estimates in both  $H^1$  and  $L^2$  norms were derived in [23] for the linear element ( $p = 1$ ) in two dimensions.

(iii) *Massjung [33] proposed an hp-unfitted discontinuous Galerkin method for Problem (1) which uses weighted averages similar as in [23] and penalizes both the jump of the solution value and that of the flux. It was proved that, for the two dimensional case, the method converges in broken  $H^1$  norm at an optimal rate with respect to  $h$  and at a suboptimal rate with respect to  $p$  by a factor of  $p$ .*

Next, we introduce the following broken (semi-)norms on the space  $V$ :

$$(14) \quad \|v\|_{1,h} := \left( \sum_{i=1}^2 \left\| a^{1/2} \nabla v \right\|_{L^2(\Omega_i)}^2 + \sum_{e \in \mathcal{I}_h} \frac{\gamma_0 p^2}{h_{K^e}} \|[v]\|_{L^2(e)}^2 + \sum_{e \in \mathcal{I}_h} \frac{\gamma_1 h_{K^e}}{p^2} \|[a \nabla v \cdot \mathbf{n}]\|_{L^2(e)}^2 \right)^{1/2},$$

$$(15) \quad \|v\|_{1,h} := \left( \|v\|_{1,h}^2 + \sum_{e \in \mathcal{I}_h} \frac{h_{K^e}}{\gamma_0 p^2} \|\{a \nabla v \cdot \mathbf{n}\}\|_{L^2(e)}^2 \right)^{1/2}.$$

### 3 Some preliminary lemmas

Recall that, for any interface segment/patch  $e \in \mathcal{I}_h$ ,  $K^e \in \mathcal{M}_h$  is an element containing  $e$  and  $K_i^e = K^e \cap \Omega_i$ ,  $i = 1, 2$ . We have the following lemma, which gives the trace and inverse trace inequalities on  $K^e$ , whose proof will be given in Section 6.

**Lemma 3.1.** *There exists a positive constant  $h_0$  depending only on the interface  $\Gamma$  and the shape regularity of the meshes, such that for all  $h \in (0, h_0]$  and any interface segment/patch  $e \in \mathcal{I}_h$ , the following estimates hold for either  $i_e = 1$  or  $i_e = 2$ :*

(i) *For any  $v \in H^1(K_{i_e}^e)$ ,*

$$(16) \quad \|v\|_{L^2(e)} \lesssim h_{K^e}^{-1/2} \|v\|_{L^2(K_{i_e}^e)} + \|v\|_{L^2(K_{i_e}^e)}^{1/2} \|\nabla v\|_{L^2(K_{i_e}^e)}^{1/2}.$$

(ii) *For any  $v_h \in \mathcal{P}_p(K^e)$ ,*

$$(17) \quad \|v_h\|_{L^2(e)} \lesssim \frac{p}{h_{K^e}^{1/2}} \|v_h\|_{L^2(K_{i_e}^e)}.$$

We remark that, if  $e \subset \partial K^e$ , then the result (i) in the above lemma is a direct consequence of the standard local trace inequality, and the estimate in (ii) is already known (cf. [10]).

The error analysis relies on the following well-known  $hp$  approximation properties (cf. [5, 20, 21]):

**Lemma 3.2.** Suppose  $w \in H^s(\Omega) \cap H_0^1(\Omega)$ ,  $s \geq 2$ . Let  $\mu = \min\{p+1, s\}$ .

(i) There exists a piecewise polynomial  $\tilde{w}_h$ ,  $\tilde{w}_h|_K \in \mathcal{P}_p(K)$ , such that, for  $j = 0, 1, \dots, s$ ,

$$\|w - \tilde{w}_h\|_{H^j(K)} \lesssim \frac{h^{\mu-j}}{p^{s-j}} \|w\|_{H^s(K)}, \quad \forall K \in \mathcal{M}_h.$$

(ii) There exists  $\hat{w}_h \in U_h^p$  such that

$$\|w - \hat{w}_h\|_{H^j(\Omega)} \lesssim \frac{h^{\mu-j}}{p^{s-j}} \|w\|_{H^s(\Omega)}, \quad j = 0, 1.$$

Here the invisible constants in the above two inequalities depend on  $s$  but are independent of  $h$  and  $p$ .  $U_h^p$  is the  $hp$ -continuous finite element space defined in (3).

The following lemma is from the standard Sobolev extension property (cf. [1]).

**Lemma 3.3.** Let  $s \geq 2$ . There exist two extension operators  $E_1 : H^s(\Omega_1) \mapsto H^s(\Omega) \cap H_0^1(\Omega)$  and  $E_2 : \{w \in H^s(\Omega_2) : w|_{\partial\Omega} = 0\} \mapsto H^s(\Omega) \cap H_0^1(\Omega)$  such that

$$(E_i w)|_{\Omega_i} = w \quad \text{and} \quad \|E_i w\|_{H^s(\Omega)} \lesssim \|w\|_{H^s(\Omega_i)}, \quad i = 1, 2,$$

where  $w \in H^s(\Omega_1)$  for  $i = 1$  and  $w \in H^s(\Omega_2)$ ,  $w|_{\partial\Omega} = 0$  for  $i = 2$  (cf. Fig. 1).

## 4 Error estimates for the symmetric interface penalty finite element method

In this section we derive the  $H^1$ - and  $L^2$ -error estimates for the  $hp$ -SIPFEM (i.e. the case  $\beta = 1$ ). The following lemma gives the continuity and coercivity of the bilinear form  $a_h(\cdot, \cdot)$  for the SIPFEM.

**Lemma 4.1.** We have

$$(18) \quad |a_h(v, w)| \leq 2 \|v\|_{1,h} \|w\|_{1,h}, \quad \forall v, w \in V.$$

For any  $0 < \gamma_1 \lesssim 1$  and  $0 < \varepsilon < 1$ , there exists a constant  $\alpha_{0,\varepsilon}$  independent of  $h$ ,  $p$ , and the penalty parameters such that, if  $\gamma_0 \geq \alpha_{0,\varepsilon}/\gamma_1$ , then, for  $0 < h \leq h_0$ ,

$$(19) \quad a_h(v_h, v_h) \geq (1 - \varepsilon) \|v_h\|_{1,h}^2, \quad \forall v_h \in V_h^p,$$

where the constant  $h_0$  is from Lemma 3.1.

*Proof.* (18) is a direct consequence of the definitions (6)–(8), (14)–(15), and the Cauchy–Schwarz inequality. It remains to prove (19). We have,

$$\begin{aligned}
(20) \quad a_h(v_h, v_h) &= \sum_{i=1}^2 \left\| a^{1/2} \nabla v_h \right\|_{L^2(\Omega_i)}^2 - 2 \sum_{e \in \mathcal{I}_h} \int_e \{a \nabla v_h \cdot \mathbf{n}\} [v_h] \\
&\quad + \sum_{e \in \mathcal{I}_h} \left( \frac{\gamma_0 p^2}{h_{K^e}} \|[v_h]\|_{L^2(e)}^2 + \frac{\gamma_1 h_{K^e}}{p^2} \|[a \nabla v_h \cdot \mathbf{n}]\|_{L^2(e)}^2 \right) \\
&= \|v_h\|_{1,h}^2 - \sum_{e \in \mathcal{I}_h} \frac{h_{K^e}}{\gamma_0 p^2} \|\{a \nabla v_h \cdot \mathbf{n}\}\|_{L^2(e)}^2 \\
&\quad - 2 \sum_{e \in \mathcal{I}_h} \int_e \{a \nabla v_h \cdot \mathbf{n}\} [v_h].
\end{aligned}$$

It is clear that, for any  $e \in \mathcal{I}_h$ ,

$$(21) \quad \{a \nabla v_h \cdot \mathbf{n}\}|_e = (a \nabla v_h)|_{K_{i_e}^e} \cdot \mathbf{n} + \frac{(-1)^{i_e}}{2} [a \nabla v_h \cdot \mathbf{n}]|_e,$$

where  $i_e$  is specified in Lemma 3.1. From Lemma 3.1 (ii), there exists a constant  $C_1$  such that,

$$(22) \quad \left\| (a \nabla v_h)|_{K_{i_e}^e} \cdot \mathbf{n} \right\|_{L^2(e)} \lesssim \frac{p}{h_{K^e}^{1/2}} \|\nabla v_h\|_{L^2(K_{i_e}^e)} \leq C_1 \frac{p}{h_{K^e}^{1/2}} \left\| a^{1/2} \nabla v_h \right\|_{L^2(K_{i_e}^e)}.$$

Therefore, from (21) and (22),

$$\begin{aligned}
(23) \quad &\sum_{e \in \mathcal{I}_h} \frac{h_{K^e}}{\gamma_0 p^2} \|\{a \nabla v_h \cdot \mathbf{n}\}\|_{L^2(e)}^2 \\
&\leq 2 \sum_{e \in \mathcal{I}_h} \frac{h_{K^e}}{\gamma_0 p^2} \left( \left\| (a \nabla v_h)|_{K_{i_e}^e} \cdot \mathbf{n} \right\|_{L^2(e)}^2 + \frac{1}{4} \|[a \nabla v_h \cdot \mathbf{n}]\|_{L^2(e)}^2 \right) \\
&\leq \sum_{e \in \mathcal{I}_h} \frac{2C_1^2}{\gamma_0} \left\| a^{1/2} \nabla v_h \right\|_{L^2(K_{i_e}^e)}^2 + \sum_{e \in \mathcal{I}_h} \frac{h_{K^e}}{2\gamma_0 p^2} \|[a \nabla v_h \cdot \mathbf{n}]\|_{L^2(e)}^2 \\
&\leq \max \left( \frac{2C_1^2}{\gamma_0}, \frac{1}{2\gamma_0 \gamma_1} \right) \|v_h\|_{1,h}^2.
\end{aligned}$$

On the other hand,

$$\begin{aligned}
(24) \quad & 2 \sum_{e \in \mathcal{I}_h} \int_e \{a \nabla v_h \cdot \mathbf{n}\} [v_h] \\
& \leq 2 \sum_{e \in \mathcal{I}_h} \left( \left\| (a \nabla v_h)|_{K_{i_e}^e} \cdot \mathbf{n} \right\|_{L^2(e)} + \frac{1}{2} \|[a \nabla v_h \cdot \mathbf{n}]\|_{L^2(e)} \right) \|[v_h]\|_{L^2(e)} \\
& \leq \sum_{e \in \mathcal{I}_h} \left( 2C_1 \left\| a^{1/2} \nabla v_h \right\|_{L^2(K_{i_e}^e)} \frac{p}{h_{K^e}^{1/2}} \|[v_h]\|_{L^2(e)} \right. \\
& \quad \left. + \|[a \nabla v_h \cdot \mathbf{n}]\|_{L^2(e)} \|[v_h]\|_{L^2(e)} \right) \\
& \leq \sum_{e \in \mathcal{I}_h} \left( \frac{4C_1^2}{\gamma_0 \varepsilon} \left\| a^{1/2} \nabla v_h \right\|_{L^2(K_{i_e}^e)}^2 + \frac{\varepsilon \gamma_0 p^2}{4 h_{K^e}} \|[v_h]\|_{L^2(e)}^2 \right. \\
& \quad \left. + \frac{1}{\gamma_0 \varepsilon} \frac{h_{K^e}}{p^2} \|[a \nabla v_h \cdot \mathbf{n}]\|_{L^2(e)}^2 + \frac{\varepsilon \gamma_0 p^2}{4 h_{K^e}} \|[v_h]\|_{L^2(e)}^2 \right) \\
& \leq \max \left( \frac{4C_1^2}{\gamma_0 \varepsilon}, \frac{1}{\gamma_0 \gamma_1 \varepsilon}, \frac{\varepsilon}{2} \right) \|v_h\|_{1,h}^2.
\end{aligned}$$

By combining (20), (23), and (24), we conclude that (19) holds. This completes the proof of the lemma.  $\square$

The following lemma is an analogue of the Cea's lemma for finite element methods.

**Lemma 4.2.** *There exists a constant  $\alpha_0$  independent of  $h$ ,  $p$ , and the penalty parameters such that for  $0 < \gamma_1 \lesssim 1$ ,  $\gamma_0 \geq \alpha_0/\gamma_1$ , and  $0 < h \leq h_0$ , the following error estimate holds.*

$$\|u - u_h\|_{1,h} \lesssim \inf_{z_h \in V_h^p} \|u - z_h\|_{1,h},$$

where the constant  $h_0$  is from Lemma 3.1.

*Proof.* From (12) and (13), we have the following Galerkin orthogonality

$$(25) \quad a_h(u - u_h, v_h) = 0, \quad \forall v_h \in V_h^p.$$

For any  $z_h \in V_h^p$ , let  $\eta_h = u_h - z_h$ . By letting  $\varepsilon = 1/2$  and in Lemma 4.1, denoting by  $\alpha_0 = \alpha_{0,1/2}$ , and using (25), we have, for  $0 < \gamma_1 \lesssim 1$  and  $\gamma_0 \geq \alpha_0/\gamma_1$ ,

$$\|\eta_h\|_{1,h}^2 \leq 2a_h(\eta_h, \eta_h) = 2a_h(u - z_h, \eta_h) \leq 4 \|u - z_h\|_{1,h} \|\eta_h\|_{1,h},$$

and hence,

$$\|u_h - z_h\|_{1,h} \leq 4 \|u - z_h\|_{1,h},$$

which implies that

$$\|u - u_h\|_{1,h} = \|u - z_h + z_h - u_h\|_{1,h} \leq 5 \|u - z_h\|_{1,h}.$$

This completes the proof of the lemma.  $\square$

The following lemma gives some approximation properties of the space  $V_h^p$ .

**Lemma 4.3.** *Let  $s \geq 2$  be an integer and let  $\mu = \min\{p+1, s\}$ . Suppose the solution to the elliptic interface problem (1) satisfies  $u|_{\Omega_i} \in H^s(\Omega_i)$ ,  $i = 1, 2$ . Then there exists  $\hat{u}_h \in V_h^p$  such that, for  $0 < h \leq h_0$ ,*

$$(26) \quad \|u - \hat{u}_h\|_{L^2(\Omega)} \lesssim \frac{h^\mu}{p^s} \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right)^{1/2},$$

$$(27) \quad \|u - \hat{u}_h\|_{1,h} \lesssim \left( \frac{1}{p} + \gamma_0 + \frac{1}{\gamma_0 p} + \frac{\gamma_1}{p} \right)^{1/2} \frac{h^{\mu-1}}{p^{s-3/2}} \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right)^{1/2},$$

where the constant  $h_0$  is from Lemma 3.1.

*Proof.* Let  $u_i = E_i u$ , where  $E_i, i = 1, 2$  are the extension operators from Lemma 3.3. Then

$$(28) \quad u_i \in H^s(\Omega) \cap H_0^1(\Omega), \quad u_i|_{\Omega_i} = u|_{\Omega_i}, \quad \|u_i\|_{H^s(\Omega)} \lesssim \|u\|_{H^s(\Omega_i)}.$$

From Lemma 3.2(ii) and (28), there exist  $\hat{u}_{ih} \in U_h^p$ ,  $i = 1, 2$  such that

$$(29) \quad \|u_i - \hat{u}_{ih}\|_{H^j(\Omega)} \lesssim \frac{h^{\mu-j}}{p^{s-j}} \|u\|_{H^s(\Omega_i)}, \quad j = 0, 1.$$

Define  $\hat{u}_h \in V_h^p$  by  $\hat{u}_h|_{\Omega_i} = \hat{u}_{ih}|_{\Omega_i}$ ,  $i = 1, 2$ . It is clear that (26) holds.

Denote by  $\zeta = u - \hat{u}_h$  and by  $\zeta_i = u_i - \hat{u}_{ih}$ ,  $i = 1, 2$ . Obviously,  $\zeta|_{\Omega_i} = \zeta_i|_{\Omega_i}$ ,  $i = 1, 2$ . Next we estimate each term in  $\|\zeta\|_{1,h}$  (cf. (14)–(15)). Clearly,

$$(30) \quad \sum_{i=1}^2 \|a^{1/2} \nabla \zeta\|_{L^2(\Omega_i)}^2 \lesssim \frac{h^{2\mu-2}}{p^{2s-2}} \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right).$$

From Lemma 3.1 (i) and (29),

$$(31) \quad \begin{aligned} & \sum_{e \in \mathcal{I}_h} \frac{\gamma_0 p^2}{h_{K^e}} \|[\zeta]\|_{L^2(e)}^2 \lesssim \sum_{e \in \mathcal{I}_h} \frac{\gamma_0 p^2}{h_{K^e}} \sum_{i=1}^2 \|\zeta_i\|_{L^2(e)}^2 \\ & \lesssim \sum_{e \in \mathcal{I}_h} \frac{\gamma_0 p^2}{h_{K^e}} \sum_{i=1}^2 \left( h_{K^e}^{-1} \|\zeta_i\|_{L^2(K^e)}^2 + \|\zeta_i\|_{L^2(K^e)} \|\nabla \zeta_i\|_{L^2(K^e)} \right) \\ & \lesssim \frac{\gamma_0 p^2}{h} \sum_{i=1}^2 \left( h^{-1} \|\zeta_i\|_{L^2(\Omega)}^2 + \|\zeta_i\|_{L^2(\Omega)} \|\nabla \zeta_i\|_{L^2(\Omega)} \right) \\ & \lesssim \gamma_0 \frac{h^{2\mu-2}}{p^{2s-3}} \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right). \end{aligned}$$

To estimate the remaining terms, we estimate  $\|(a \nabla \zeta)|_{\Omega_i} \cdot \mathbf{n}\|_{L^2(e)}^2$ ,  $i = 1, 2, e \in \mathcal{I}_h$ . From Lemma 3.2 (i), there exist piecewise polynomials  $\tilde{u}_{ih}, \tilde{u}_{ih}|_K \in \mathcal{P}_p(K)$ ,

such that

$$\|u_i - \tilde{u}_{ih}\|_{H^j(K)} \lesssim \frac{h^{\mu-j}}{p^{s-j}} \|u_i\|_{H^s(K)}, \quad \forall j = 0, 1, 2, \quad i = 1, 2, \quad K \in \mathcal{M}_h.$$

Therefore from Lemma 3.1,

$$\begin{aligned} \|(a\nabla\zeta)|_{\Omega_i} \cdot \mathbf{n}\|_{L^2(e)}^2 &\lesssim \|\nabla(u_i - \tilde{u}_{ih})\|_{L^2(e)}^2 + \|\nabla(\tilde{u}_{ih} - \hat{u}_{ih})\|_{L^2(e)}^2 \\ &\lesssim h_{K^e}^{-1} \|\nabla(u_i - \tilde{u}_{ih})\|_{L^2(K^e)}^2 \\ &\quad + \|\nabla(u_i - \tilde{u}_{ih})\|_{L^2(K^e)} \|u_i - \tilde{u}_{ih}\|_{H^2(K^e)} \\ &\quad + \frac{p^2}{h_{K^e}} \|\nabla(\tilde{u}_{ih} - u_i + u_i - \hat{u}_{ih})\|_{L^2(K^e)}^2 \\ &\lesssim \frac{h^{2\mu-3}}{p^{2s-4}} \|u_i\|_{H^s(K^e)}^2 + \frac{p^2}{h} \|\nabla(u_i - \hat{u}_{ih})\|_{L^2(K^e)}^2. \end{aligned}$$

It follows from (28) and (29) that

$$\sum_{e \in \mathcal{I}_h} \|(a\nabla\zeta)|_{\Omega_i} \cdot \mathbf{n}\|_{L^2(e)}^2 \lesssim \frac{h^{2\mu-3}}{p^{2s-4}} \|u\|_{H^s(\Omega_i)}^2.$$

Thus,

$$\begin{aligned} (32) \quad \sum_{e \in \mathcal{I}_h} \left( \frac{\gamma_1 h_{K^e}}{p^2} \|[a\nabla\zeta \cdot \mathbf{n}]\|_{L^2(e)}^2 + \frac{h_{K^e}}{\gamma_0 p^2} \|\{a\nabla\zeta \cdot \mathbf{n}\}\|_{L^2(e)}^2 \right) \\ \lesssim \left( \gamma_1 + \frac{1}{\gamma_0} \right) \frac{h^{2\mu-2}}{p^{2s-2}} \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right) \end{aligned}$$

Then (27) follows by combining (30)–(32). This completes the proof of the lemma.  $\square$

By combining Lemma 4.2 and Lemma 4.3, we have the following theorem which gives error estimate in the norm  $\|\cdot\|_{1,h}$  for the SIPFEM. The proof is straightforward and is omitted.

**Theorem 4.4.** *Let  $s \geq 2$  be an integer and let  $\mu = \min\{p+1, s\}$ . Suppose the solution to the elliptic interface problem (1) satisfies  $u|_{\Omega_i} \in H^s(\Omega_i)$ ,  $i = 1, 2$ . Then there exists a constant  $\alpha_0$  independent of  $h, p$ , and the penalty parameters such that for  $0 < \gamma_1 \lesssim 1$  and  $\gamma_0 \geq \alpha_0/\gamma_1$ , the following error estimate holds:*

$$\|u - u_h\|_{1,h} \lesssim \gamma_0^{1/2} \frac{h^{\mu-1}}{p^{s-3/2}} \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right)^{1/2}, \quad \forall 0 < h \leq h_0,$$

where the constant  $h_0$  from Lemma 3.1 depends only on the interface  $\Gamma$  and the shape regularity of the meshes.

Next we derive the  $L^2$ -error estimate by using the Nitsche's duality argument (cf. [15]). Consider the following auxiliary problem:

$$(33) \quad \begin{cases} -\nabla \cdot (a(x)\nabla w) = u - u_h, & \text{in } \Omega_1 \cup \Omega_2, \\ [w] = 0, \quad [(a(x)\nabla w) \cdot \mathbf{n}] = 0, & \text{on } \Gamma, \\ w = 0, & \text{on } \partial\Omega, \end{cases}$$

It can be shown that  $w$  satisfies (cf. [4])

$$(34) \quad \|w\|_{H^2(\Omega_1)} + \|w\|_{H^2(\Omega_2)} \lesssim \|u - u_h\|_{L^2(\Omega)}.$$

**Theorem 4.5.** *Under the conditions of Theorem 4.4, there holds the following estimate for the SIPFEM.*

$$\|u - u_h\|_{L^2(\Omega)} \lesssim \gamma_0 \frac{h^\mu}{p^{s-1}} \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right)^{1/2}, \quad \forall 0 < h \leq h_0,$$

where the constant  $h_0$  is from Lemma 3.1.

*Proof.* Let  $\eta = u - u_h$ . Testing (33) by  $\eta$  and using (25) we get

$$(35) \quad \|\eta\|_{L^2(\Omega)}^2 = a_h(w, \eta) = a_h(\eta, w) = a_h(\eta, w - \hat{w}_h),$$

where  $\hat{w}_h \in V_h^p$  satisfies the estimate (cf. Lemma 4.3)

$$(36) \quad \|w - \hat{w}_h\|_{1,h} \lesssim \gamma_0^{1/2} \frac{h}{p^{1/2}} \sum_{i=1}^2 \|w\|_{H^2(\Omega_i)}.$$

Therefore from (18) and (34),

$$\|\eta\|_{L^2(\Omega)}^2 \leq 2 \|w - \hat{w}_h\|_{1,h} \|\eta\|_{1,h} \lesssim \gamma_0^{1/2} \frac{h}{p^{1/2}} \|\eta\|_{L^2(\Omega)} \|\eta\|_{1,h},$$

that is  $\|\eta\|_{L^2(\Omega)} \lesssim \gamma_0^{1/2} \frac{h}{p^{1/2}} \|\eta\|_{1,h}$ , which completes the proof of Theorem 4.5 by using Theorem 4.4.  $\square$

## 5 Error estimates for the non-symmetric interface penalty finite element methods

In this section we derive the  $H^1$ - and  $L^2$ - error estimates for the  $hp$ -NIPFEM (i.e. the case  $\beta = -1$ ). The following lemma gives the continuity and coercivity of the bilinear form  $a_h(\cdot, \cdot)$  for the NIPFEM.

**Lemma 5.1.** *We have, for the NIPFEM ( $\beta = -1$ ),*

$$(37) \quad |a_h(v, w)| \leq 2 \|v\|_{1,h} \|w\|_{1,h}, \quad \forall v, w \in V.$$

$$(38) \quad a_h(v_h, v_h) \gtrsim \frac{\gamma_0 \gamma_1}{\gamma_0 \gamma_1 + \max(\gamma_1, 1)} \|v_h\|_{1,h}^2, \quad \forall v_h \in V_h^p, \quad \forall 0 < h \leq h_0,$$

where the constant  $h_0$  is from Lemma 3.1.

*Proof.* (37) is a direct consequence of the definitions (6)–(8), (14)–(15), and the Cauchy–Schwarz inequality. From (23),

$$\|\|v_h\|\|_{1,h}^2 \leq \left(1 + \max\left(\frac{2C_1^2}{\gamma_0}, \frac{1}{2\gamma_0\gamma_1}\right)\right) \|v_h\|_{1,h}^2.$$

Then (38) follows from the identity  $a_h(v_h, v_h) = \|v_h\|_{1,h}^2$ . This completes the proof of the lemma.  $\square$

Following the argument of Theorem 4.4 we have the following error estimate in the norm  $\|\|\cdot\|\|_{1,h}$  for the *hp*-NIPFEM. The proof is omitted.

**Theorem 5.2.** *Let  $s \geq 2$  be an integer and let  $\mu = \min\{p+1, s\}$ . Suppose the solution to the elliptic interface problem (1) satisfies  $u|_{\Omega_i} \in H^s(\Omega_i)$ ,  $i = 1, 2$ . For  $0 < \gamma_1 \lesssim 1$  and  $\gamma_0 \gtrsim 1/\gamma_1$ , the following error estimate holds.*

$$\|\|u - u_h\|\|_{1,h} \lesssim \gamma_0^{1/2} \frac{h^{\mu-1}}{p^{s-3/2}} \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right)^{1/2}, \quad \forall 0 < h \leq h_0,$$

where the constant  $h_0$  is from Lemma 3.1.

The following theorem gives the  $L^2$ -error estimate for the *hp*-NIPFEM.

**Theorem 5.3.** *Under the conditions of Theorem 5.2, there holds the following estimate for the NIPFEM.*

$$\|u - u_h\|_{L^2(\Omega)} \lesssim \left( \gamma_0 \frac{h^\mu}{p^{s-1}} + \frac{h^{\mu-1/2}}{p^{s-1/2}} \right) \left( \sum_{i=1}^2 \|u\|_{H^s(\Omega_i)}^2 \right)^{1/2}, \quad \forall 0 < h \leq h_0,$$

where the constant  $h_0$  is from Lemma 3.1.

*Proof.* The proof is similar to that of Theorem 4.5. So we only sketch the differences here. (35) becomes

$$\begin{aligned} (39) \quad \|\eta\|_{L^2(\Omega)}^2 &= a_h(w, \eta) = a_h(\eta, w) - 2 \sum_{e \in \mathcal{I}_h} \int_e [\eta] \{a \nabla w \cdot \mathbf{n}\} \\ &= a_h(\eta, w - \hat{w}_h) - 2 \sum_{e \in \mathcal{I}_h} \int_e [\eta] (a \nabla w \cdot \mathbf{n}), \end{aligned}$$

where  $\eta = u - u_h$ ,  $w$  is defined in (33), and  $\hat{w}_h \in V_h^p$  satisfies the estimate in (36). From (37), (34), and (36),

$$(40) \quad a_h(\eta, w - \hat{w}_h) \lesssim \gamma_0^{1/2} \frac{h}{p^{1/2}} \|\eta\|_{L^2(\Omega)} \|\| \eta \| \|_{1,h},$$

The last term in (39) is bounded by using the Cauchy-Schwarz inequality and the trace inequality:

$$\begin{aligned}
(41) \quad -2 \sum_{e \in \mathcal{I}_h} \int_e [\eta] (a \nabla w \cdot \mathbf{n}) &\leq 2 \sum_{e \in \mathcal{I}_h} \|[\eta]\|_{L^2(e)} \|a \nabla w \cdot \mathbf{n}\|_{L^2(e)} \\
&\lesssim J_0(\eta, \eta)^{1/2} \left( \sum_{e \in \mathcal{I}_h} \frac{h_{K^e}}{\gamma_0 p^2} \|\nabla w\|_{L^2(e)}^2 \right)^{1/2} \\
&\lesssim \frac{h^{1/2}}{\gamma_0^{1/2} p} J_0(\eta, \eta)^{1/2} \|\nabla w\|_{L^2(\Gamma)} \\
&\lesssim \frac{h^{1/2}}{\gamma_0^{1/2} p} J_0(\eta, \eta)^{1/2} \left( \|\nabla w\|_{L^2(\Omega_1)} \|\nabla w\|_{H^1(\Omega_1)} \right)^{1/2} \\
&\lesssim \frac{h^{1/2}}{\gamma_0^{1/2} p} \|\eta\|_{1,h} \|\eta\|_{L^2(\Omega)},
\end{aligned}$$

where we have used (34) to derive the last inequality. By combining (39)–(41) we conclude that

$$\|\eta\|_{L^2(\Omega)} \lesssim \left( \gamma_0^{1/2} \frac{h}{p^{1/2}} + \frac{h^{1/2}}{\gamma_0^{1/2} p} \right) \|\eta\|_{1,h},$$

which completes the proof of Theorem 5.3 by using Theorem 5.2.  $\square$

**Remark 5.1.** *Here is a brief comparison between the SIPFEM and the NIPFEM:*

- The stiffness matrix of the SIPFEM is symmetric while that of the NIPFEM is not.
- Both methods require that the penalty parameters satisfy  $\gamma_0 \gamma_1 \geq C$  but  $C$  is a problem dependent constant for the SIPFEM while  $C$  can be an arbitrary positive constant for the NIPFEM.
- The  $L^2$ -error estimate for the SIPFEM is optimal with respect to  $h$  while that for the NIPFEM is not.

## 6 Proof of Lemma 3.1

In this section we prove the local trace inequality and the local inverse trace inequality in Lemma 3.1. The proof is divided into two cases, the two dimensional case and the three dimensional case. Recall that,  $e \in \mathcal{I}_h$  is an interface segment/patch,  $K^e \in \mathcal{M}_h$  is an element containing  $e$ , and  $K_i^e = K^e \cap \Omega_i$ ,  $i = 1, 2$ .

### 6.1 Two dimensional case

Since  $\Gamma$  is a  $C^2$  boundary of  $\Omega_1$ , it can be expressed as a union of open sub-segments  $\Gamma_i, i \in \Lambda := \{1, 2, \dots, N\}$  such that each  $\Gamma_i$  is parameterized by some

function  $x = \mathbf{r}_i(\xi) \in [C^2(I_i)]^2$  satisfying  $|\mathbf{r}'_i(\xi)| \neq 0$ , where  $I_i$  is an open interval. Since  $\Gamma_i \cap \Gamma_j$  is either empty or an open subsegment of  $\Gamma$ ,  $e$  is contained entirely in some  $\Gamma_k$  if  $h_{K^e}$  is small enough, and hence  $e$  is parameterized by  $x = \mathbf{r}(\xi), \xi \in I$  with  $\mathbf{r}(\xi) = \mathbf{r}_k(\xi)$  and  $I \subset I_k$ . It is clear that the length of  $I \lesssim h_{K^e}$ .

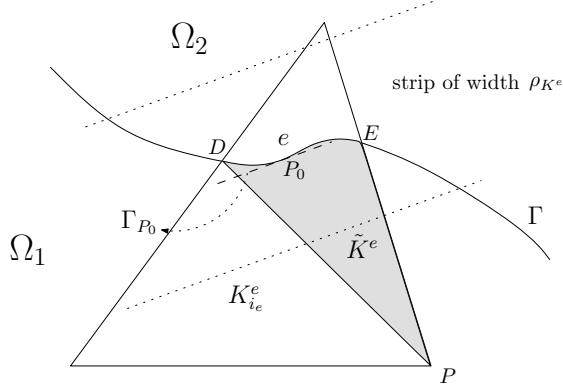


Figure 2: An interface segment  $e$  and the element  $K^e$  in 2D

Let  $P_0$  be any fixed point on  $e$  and let  $\Gamma_{P_0}$  be the tangent line to  $\Gamma$  at  $P_0$ . Let  $P$  be a point in  $K^e$  which achieves the maximum distance, denoted by  $\tilde{h}$ , from the tangent line  $\Gamma_{P_0}$ . Denote by  $\rho_{K^e}$  the diameter of the largest disk contained in  $K^e$ . Clearly, we have  $\tilde{h} \geq \rho_{K^e}/2 \gtrsim h_{K^e}$  (see Fig. 2, otherwise  $K^e$  would be completely contained in a strip of width less than  $\rho_{K^e}$  along the axis  $\Gamma_{P_0}$ , which causes a contradiction). Now, we may choose  $K_{i_e}^e$  to be the one of two subregions  $K_1^e$  and  $K_2^e$  which contains the point  $P$  (in Fig. 2, for instance,  $K_{i_e}^e = K_1^e$ ). Without loss of generality, we assume  $P$  is the origin. Suppose  $P_0$  corresponds to the parameter  $\xi = \xi_0$ . Then the vector  $\overrightarrow{PP_0} = \mathbf{r}(\xi_0)$ . It is clear that

$$(42) \quad |\mathbf{r}(\xi_0)| \leq h_{K^e}, \quad |\mathbf{r}'(\xi_0)| \approx 1.$$

We have the following Taylor's expansions at  $\xi_0$ ,

$$(43) \quad \mathbf{r}(\xi) = \mathbf{r}(\xi_0) + \mathbf{r}'(\xi_0)(\xi - \xi_0) + O(h_{K^e}^2),$$

$$(44) \quad \mathbf{r}'(\xi) = \mathbf{r}'(\xi_0) + O(h_{K^e}).$$

To proceed, we need to bound the following term  $G(\xi)$  from below:

$$G(\xi) := \frac{|\mathbf{r}(\xi) \times \mathbf{r}'(\xi)|}{|\mathbf{r}'(\xi)|},$$

where  $\mathbf{a} \times \mathbf{b} = a_1b_2 - a_2b_1$  denotes the scalar cross product of two dimensional vectors. Recall that  $|\mathbf{a} \times \mathbf{b}| = |\mathbf{a}||\mathbf{b}| \sin\langle \mathbf{a}, \mathbf{b} \rangle$ . Note that  $G(\xi_0) = \frac{|\mathbf{r}(\xi_0) \times \mathbf{r}'(\xi_0)|}{|\mathbf{r}'(\xi_0)|}$

equals to the distance from  $P$  to  $\Gamma_{P_0}$ , thus

$$G(\xi_0) = \tilde{h} \gtrsim h_{K^e}.$$

By combining (42)–(44), we derive that

$$G(\xi) = \frac{|\mathbf{r}(\xi_0) \times \mathbf{r}'(\xi_0) + O(h_{K^e}^2)|}{|\mathbf{r}'(\xi_0) + O(h_{K^e})|} = G(\xi_0) + O(h_{K^e}^2).$$

Then there exists a positive constant  $h_0$  depending only on the interface  $\Gamma$  and the shape regularity of  $K^e$ , such that if  $0 < h_{K^e} \leq h_0$ , then

$$(45) \quad G(\xi) \gtrsim h_{K^e}, \quad \text{for all } \xi \in I.$$

Next we prove the trace inequality (16). Let  $\tilde{K}^e$  be a fan-like region in  $K_{i_e}^e$ , with  $e \cup PD \cup PE$  as its boundary (cf. Fig. 2), that is,

$$\tilde{K}^e := \{x : x = t \cdot \mathbf{r}(\xi), t \in [0, 1], \xi \in I\}.$$

Let  $v \in C^1(K_{i_e}^e)$  and consider its restriction on  $e$  as follows,

$$(46) \quad v^2(\mathbf{r}(\xi)) = \int_0^1 \frac{\partial}{\partial t} ((t^2 v^2(t \cdot \mathbf{r}(\xi))) \, dt = \int_0^1 2(t^2 v v_t + t v^2) \, dt.$$

From  $v_t = \nabla v \cdot \mathbf{r}$ , we have

$$v^2(\mathbf{r}(\xi)) \leq \int_0^1 2(|v| |\nabla v| |\mathbf{r}| + v^2) t \, dt.$$

Integrating along  $e$ , we find

$$\begin{aligned} (47) \quad \int_e v^2 \, ds &= \int_I v^2(\mathbf{r}(\xi)) |\mathbf{r}'(\xi)| \, d\xi \\ &\leq 2 \int_I \int_0^1 (|v| |\nabla v| |\mathbf{r}| + v^2) t |\mathbf{r}'(\xi)| \, dt \, d\xi \\ &= 2 \int_I \int_0^1 (|v| |\nabla v| |\mathbf{r}| + v^2) \cdot t |\mathbf{r}(\xi) \times \mathbf{r}'(\xi)| \cdot \frac{1}{G(\xi)} \, dt \, d\xi \\ &\lesssim \frac{1}{\inf_{\xi \in I} G(\xi)} \int_I \int_0^1 (|v| |\nabla v| h_{K^e} + v^2) \cdot t |\mathbf{r}(\xi) \times \mathbf{r}'(\xi)| \, dt \, d\xi, \end{aligned}$$

where we have used  $|\mathbf{r}(\xi)| \lesssim h_{K^e}$  to derive the last inequality. By noting that  $t |\mathbf{r}(\xi) \times \mathbf{r}'(\xi)|$  is the absolute value of the Jacobian determinant of the mapping  $x = t \cdot \mathbf{r}(\xi)$  and using (45), we obtain that

$$\begin{aligned} (48) \quad \|v\|_{L^2(e)}^2 &\lesssim \frac{1}{\inf_{\xi \in I} G(\xi)} \int_{\tilde{K}^e} (|v| |\nabla v| h_{K^e} + v^2) \, dx \\ &\lesssim \|v\|_{L^2(\tilde{K}^e)} \|\nabla v\|_{L^2(\tilde{K}^e)} + h_{K^e}^{-1} \|v\|_{L^2(\tilde{K}^e)}^2. \end{aligned}$$

Since  $\tilde{K}^e \subseteq K_{i_e}^e$ , we conclude that

$$\|v\|_{L^2(e)} \lesssim h_{K^e}^{-\frac{1}{2}} \|v\|_{L^2(K_{i_e}^e)} + \|v\|_{L^2(K_{i_e}^e)}^{\frac{1}{2}} \|\nabla v\|_{L^2(K_{i_e}^e)}^{\frac{1}{2}}.$$

By a density argument, the above inequality is also valid for  $v \in H^1(K_{i_e}^e)$ . This completes the proof of (16).

It remains to prove the inverse trace inequality (17). We need the following inverse inequality on  $\mathcal{P}_p([0, 1])$  (cf. [36]):

$$(49) \quad \|w'_h\|_{L^2([0,1])} \lesssim p^2 \|w_h\|_{L^2([0,1])} \quad \forall w_h \in \mathcal{P}_p([0, 1]).$$

By noting that  $v_h(t \cdot \mathbf{r}(\xi))$  is a polynomial in  $t$  of degree  $\leq p$  if  $K^e$  is a triangle and of degree  $\leq 2p$  if  $K^e$  is a parallelogram, we have from the above inverse inequality that

$$(50) \quad \begin{aligned} v_h^2(\mathbf{r}(\xi)) &= \int_0^1 \frac{\partial}{\partial t} (t^2 v_h^2(t \cdot \mathbf{r}(\xi))) dt = 2 \int_0^1 t v_h(t \cdot \mathbf{r}(\xi)) \frac{\partial}{\partial t} (t v_h(t \cdot \mathbf{r}(\xi))) dt \\ &\leq 2 \left( \int_0^1 (t v_h(t \cdot \mathbf{r}(\xi)))^2 dt \right)^{\frac{1}{2}} \left( \int_0^1 \left( \frac{\partial}{\partial t} (t v_h(t \cdot \mathbf{r}(\xi))) \right)^2 dt \right)^{\frac{1}{2}} \\ &\lesssim p^2 \int_0^1 t^2 v_h^2(t \cdot \mathbf{r}(\xi)) dt. \end{aligned}$$

Following the same guidelines of (47) and (48), we derive that

$$(51) \quad \begin{aligned} \|v_h\|_{L^2(e)}^2 &= \int_I v_h^2(\mathbf{r}(\xi)) |\mathbf{r}'(\xi)| d\xi \lesssim p^2 \int_I \int_0^1 v_h^2(t \cdot \mathbf{r}(\xi)) t |\mathbf{r}'(\xi)| dt d\xi \\ &\lesssim \frac{p^2}{\inf_{\xi \in I} G(\xi)} \int_{\tilde{K}^e} v_h^2 dx \lesssim \frac{p^2}{h_{K^e}} \|v_h\|_{L^2(\tilde{K}^e)}^2, \end{aligned}$$

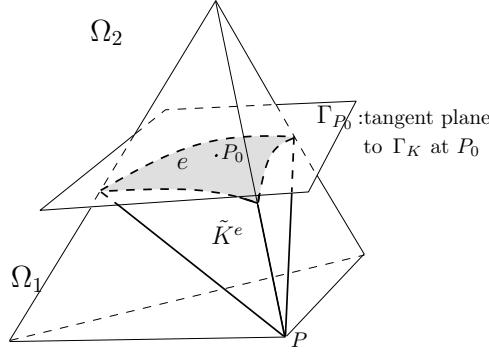
which yields the conclusion of (17).

## 6.2 Three dimensional case

The argument for the two dimensional case is readily extended to the three dimensional one. We only sketch the proof by indicating the necessary modifications in the proof for the two dimensional case.

Since  $\Gamma$  is a  $C^2$  interface, it can be expressed as a union of open subpatches  $\Gamma_i, i \in \Lambda := \{1, 2, \dots, N\}$  such that each  $\Gamma_i$  is parameterized by some function  $x = \mathbf{r}_i(\xi, \eta) \in [C^2(U_i)]^3$  satisfying  $|\mathbf{r}_{i\xi} \times \mathbf{r}_{i\eta}| \neq 0$ , where  $U_i$  is an open domain in  $\mathbb{R}^2$ . Since  $\Gamma_i \cap \Gamma_j$  is either empty or an open subpatch of  $\Gamma$ ,  $e$  is contained entirely in some  $\Gamma_k$  if  $h_{K^e}$  is small enough, and hence  $e$  is parameterized by  $x = \mathbf{r}(\xi, \eta), (\xi, \eta) \in U$  with  $\mathbf{r}(\xi, \eta) = \mathbf{r}_k(\xi, \eta)$  and  $U \subset U_k$ . It is clear that the diameter of  $U \lesssim h_{K^e}$ .

Let  $P_0$  be any fixed point on  $e$  and let  $\Gamma_{P_0}$  be the tangent plane to  $\Gamma$  at  $P_0$ . Then there exists a point  $P \in K^e$  such that the distance from  $P$  to  $\Gamma_{P_0} \gtrsim h_{K^e}$ .

Figure 3: An interface patch  $e$  and the element  $K^e$  in 3D

Without loss of generality, we assume  $P$  is the origin. Suppose  $P_0$  corresponds to the parameter  $(\xi, \eta) = (\xi_0, \eta_0)$ . We have

$$(52) \quad |\mathbf{r}(\xi_0, \eta_0)| \leq h_{K^e}, \quad |\mathbf{r}_\xi(\xi_0, \eta_0) \times \mathbf{r}_\eta(\xi_0, \eta_0)| \approx 1.$$

$$(53) \quad \mathbf{r}(\xi, \eta) = \mathbf{r}(\xi_0, \eta_0) + \mathbf{r}_\xi(\xi_0, \eta_0)(\xi - \xi_0) + \mathbf{r}_\eta(\xi_0, \eta_0)(\eta - \eta_0) + O(h_{K^e}^2),$$

$$(54) \quad \mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta) = \mathbf{r}_\xi(\xi_0, \eta_0) \times \mathbf{r}_\eta(\xi_0, \eta_0) + O(h_{K^e}).$$

Let

$$G(\xi, \eta) := \frac{|\mathbf{r}(\xi, \eta) \cdot (\mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta))|}{|\mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta)|}.$$

By noting that  $G(\xi_0, \eta_0)$  is the distance from  $P$  to  $\Gamma_{P_0}$ , we have  $G(\xi_0, \eta_0) \gtrsim h_{K^e}$ . From (52)–(54), there exists a positive constant  $h_0$  depending only on the interface  $\Gamma$  and the shape regularity of  $K^e$ , such that if  $0 < h_{K^e} \leq h_0$ , then

$$(55) \quad G(\xi, \eta) \gtrsim h_{K^e} \quad \forall (\xi, \eta) \in U.$$

Next we prove the trace inequality (16). Let  $\tilde{K}^e$  be the cone with apex  $P$  and curved base  $e$  as shown in Fig. 3, i.e.,

$$\tilde{K}^e := \{x : x = t \cdot \mathbf{r}(\xi, \eta), t \in [0, 1], (\xi, \eta) \in U\}.$$

We have

$$(56) \quad v^2(\mathbf{r}(\xi, \eta)) = \int_0^1 \frac{\partial}{\partial t} (t^3 v^2(t \cdot \mathbf{r}(\xi, \eta))) dt = \int_0^1 (2t^3 v v_t + 3t^2 v^2) dt.$$

Therefore, from  $|\mathbf{r}| \lesssim h_{K^e}$  and the fact that  $t^2 |\mathbf{r}(\xi, \eta) \cdot (\mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta))|$  is the absolute value of the Jacobian determinant of the mapping  $x = t \cdot \mathbf{r}(\xi, \eta)$ ,

we conclude that

$$\begin{aligned}
(57) \quad & \int_e v^2 \, d\sigma = \int_U v^2(\mathbf{r}(\xi, \eta)) |\mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta)| \, d\xi \, d\eta \\
& \lesssim \int_U \int_0^1 (|v| |\nabla v| |\mathbf{r}| + v^2) t^2 |\mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta)| \, dt \, d\xi \, d\eta \\
& \lesssim \int_U \int_0^1 (|v| |\nabla v| h_{K^e} + v^2) t^2 |\mathbf{r}(\xi, \eta) \cdot (\mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta))| \frac{1}{G(\xi, \eta)} \, dt \, d\xi \, d\eta \\
& \lesssim \frac{1}{\inf_{(\xi, \eta) \in U} G(\xi, \eta)} \int_{\tilde{K}^e} (|v| |\nabla v| h_{K^e} + v^2) \, dx \\
& \lesssim \|v\|_{L^2(\tilde{K}^e)} \|\nabla v\|_{L^2(\tilde{K}^e)} + h_{K^e}^{-1} \|v\|_{L^2(\tilde{K}^e)}^2
\end{aligned}$$

which implies that (16) holds.

It remains to prove (17). By following the proof of (50), we have

$$v_h^2(\mathbf{r}(\xi, \eta)) \lesssim p^2 \int_0^1 t^2 v_h^2(t \cdot \mathbf{r}(\xi, \eta)) \, dt.$$

Then, from (55),

$$\begin{aligned}
(58) \quad & \int_e v_h^2 \, d\sigma = \int_U v_h^2(\mathbf{r}(\xi, \eta)) |\mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta)| \, d\xi \, d\eta \\
& \lesssim p^2 \int_U \int_0^1 v_h^2(t \cdot \mathbf{r}(\xi, \eta)) t^2 |\mathbf{r}_\xi(\xi, \eta) \times \mathbf{r}_\eta(\xi, \eta)| \, dt \, d\xi \, d\eta \\
& \lesssim \frac{p^2}{\inf_{(\xi, \eta) \in U} G(\xi, \eta)} \int_{\tilde{K}^e} v_h^2 \, dx \lesssim \frac{p^2}{h_{K^e}} \|v_h\|_{L^2(\tilde{K}^e)}^2,
\end{aligned}$$

which implies that (17) holds. This completes the proof of Lemma 3.1.

## References

- [1] R. ADAMS, *Sobolev Spaces*, Academic Press, 1975.
- [2] D. ARNOLD, *An interior penalty finite element method with discontinuous elements*, SIAM J. Numer. Anal., 19 (1982), pp. 742–760.
- [3] D. ARNOLD, F. BREZZI, B. COCKBURN, AND D. MARINI, *Unified analysis of discontinuous Galerkin methods for elliptic problems.*, SIAM J. Numer. Anal., 39 (2001), pp. 1749–1779.
- [4] I. BABUŠKA, *The finite element method for elliptic equations with discontinuous coefficients*, Computing, 5 (1970), pp. 207–218.
- [5] I. BABUŠKA AND M. SURI, *The  $h$ - $p$  version of the finite element method with quasiuniform meshes.*, Math. Modeling Numer. Anal., 21 (1987), pp. 199–238.

- [6] G.A. BAKER, *Finite element methods for elliptic equations using nonconforming elements*, Math. Comp., 31 (1977), pp. 44–59.
- [7] J. W. BARRETT AND C. M. ELLIOTT, *Fitted and unfitted finite-element methods for elliptic equations with smooth interfaces*, IMA J. Numer. Anal., 7 (1987), pp. 283–300.
- [8] P. BASTIAN AND C. ENGWER, *An unfitted finite element method using discontinuous Galerkin*, Int. J. Numer. Meth. Engng, 79 (2009), pp. 1557–1576.
- [9] S. C. BRENNER AND L. R. SCOTT, *The mathematical theory of finite element methods*, Springer-Verlag, third ed., 2008.
- [10] E. BURMAN AND A. ERN, *Continuous interior penalty hp-finite element methods for advection and advection-diffusion equations*, Math. Comp., 259 (2007), pp. 1119–1140.
- [11] Z. CHEN, Y. XIAO, AND L. ZHANG, *The adaptive immersed interface finite element method for elliptic and Maxwell interface problems*, J. Comput. Phys., 228 (2009), pp. 5000–5019.
- [12] Z. CHEN AND J. ZOU, *Finite element methods and their convergence for elliptic and parabolic interface problems*, Numer. Math., 79 (1998), pp. 175–202.
- [13] S. CHOU, D. KWAK, AND K. WEE, *Optimal convergence analysis of an immersed interface finite element method*, Advances in Computational Mathematics, 33 (2010), pp. 149–168.
- [14] C. CHU, I.G. GRAHAM, AND T. HOU, *A new multiscale finite element method for high-contrast elliptic interface problems*, Math. Comp., 79 (2010), pp. 1915–1955.
- [15] P. G. CIARLET, *The Finite Element Method for Elliptic Problems*, North-Holland, Amsterdam, 1978.
- [16] J. DOUGLAS AND T. DUPONT, *Interior Penalty Procedures for Elliptic and Parabolic Galerkin methods*, Lecture Notes in Phys. 58, Springer-Verlag, Berlin, 1976.
- [17] R. P. FEDKIW, T. ASLAM, B. MERRIMAN, AND S. OSHER, *A non-oscillatory Eulerian approach to interfaces in multimaterial flows (the ghost fluid method)*, J. Comput. Phys., 152 (1999), pp. 457–492.
- [18] X. FENG AND H. WU, *Discontinuous Galerkin methods for the Helmholtz equation with large wave numbers.*, SIAM J. Numer. Anal., 47 (2009), pp. 2872–2896, also downloadable at <http://arXiv.org/abs/0810.1475>.

- [19] Y. GONG, B. LI, AND Z. LI, *Immersed-interface finite-element methods for elliptic interface problems with nonhomogeneous jump conditions*, SIAM J. Numer. Anal., 46 (2008), pp. 472–495.
- [20] B. GUO, *Approximation theory for the  $p$ -version of the finite element method in three dimensions, part 1: approximabilities of singular functions in the framework of the Jacobi-weighted Besov and Sobolev spaces.*, SIAM J. Numer. Anal., 44 (2006), pp. 246–269.
- [21] B. GUO AND W. SUN, *The optimal convergence of the  $h$ - $p$  version of the finite element method with quasi-uniform meshes.*, SIAM J. Numer. Anal., 45 (2007), pp. 698–730.
- [22] G. GUYOMARC'H, C.O. LEE, AND K. JEON, *A discontinuous Galerkin method for elliptic interface problems with application to electroporation*, Commun. Numer. Meth. Engng., 25 (2009), pp. 991–1008.
- [23] A. HANSBO AND P. HANSBO, *An unfitted finite element method, based on Nitsche's method, for elliptic interface problems*, Comput. Methods Appl. Mech. Engrg., 191 (2002), pp. 5537–5552.
- [24] T. HOU, Z. LI, S. OSHER, AND H. ZHAO, *A hybrid method for moving interface problems with application to the HeleShaw flow*, J. Comput. Phys., 134 (1997), pp. 236–252.
- [25] H. HUANG AND Z. LI, *Convergence analysis of the immersed interface method*, IMA J. Numer. Anal., 19 (1999), pp. 583–608.
- [26] R. LEVEQUE AND Z. LI, *The immersed interface method for elliptic equations with discontinuous coefficients and singular sources*, SIAM J. Numer. Anal., 31 (1994), pp. 1019–1044.
- [27] J. LI, J. M. MELENK, B. WOHLMUTH, AND J. ZOU, *Optimal a priori estimates for higher order finite elements for elliptic interface problems*, Appl. Numer. Math., 60 (2010), pp. 19–37.
- [28] Z. LI, T. LIN, AND X. WU, *New cartesian grid methods for interface problems using the finite element formulation*, Numer. Math., 96 (2003), pp. 61–98.
- [29] Z. LI, W. WANG, I.L. CHERN, AND M. LAI, *New formulations for interface problems in polar coordinates*, SIAM J. Sci. Comput., 25 (2003), pp. 224–245.
- [30] S. LIM AND C.S. PESKIN, *Simulations of the whirling instability by the immersed boundary method*, SIAM J. Sci. Comput., 25 (2004), pp. 2066–2083.
- [31] M.N. LINNICK AND H.F. FAESL, *A high-order immersed interface method for simulating unsteady incompressible flows on irregular domains*, J. Comput. Phys., 204 (2005), p. 157192.

- [32] X. LIU, R.P. FEDKIW, AND M. KANG, *A boundary condition capturing method for Poisson's equation on irregular domains*, J. Comput. Phys., 160 (2000), pp. 151–178.
- [33] R. MASSJUNG, *An hp-error estimate for an unfitted discontinuous Galerkin method applied to elliptic interface problems*, RWTH 300, IGPM Report, 2009.
- [34] C.S. PESKIN, *Numerical analysis of blood flow in heart*, J. Comput. Phys., 25 (1977), pp. 220–252.
- [35] B. RIVIÈRE, M. F. WHEELER, AND V. GIRAUT, *Improved energy estimates for interior penalty, constrained and discontinuous Galerkin methods for elliptic problems. I*, Comput. Geosci., 3 (1999), pp. 337–360.
- [36] C. SCHWAB, *P- and hp-Finite Element Methods*, Oxford University Press, 1998.
- [37] J.R. SHEWCHUK, *What is a good linear finite element? interpolation, conditioning, anisotropy, and quality measures (preprint)*, University of California at Berkeley, (2002).
- [38] Y.H. TSENG AND J.H. FERZIGER, *A ghost-cell immersed boundary method for flow in complex geometry*, J. Comput. Phys., 192 (2003), pp. 593–623.
- [39] M. F. WHEELER, *An elliptic collocation-finite element method with interior penalties*, SIAM J. Numer. Anal., 15 (1978), pp. 152–161.
- [40] S. YU, Y. ZHOU, AND G. WEI, *Matched interface and boundary (MIB) method for elliptic problems with sharp-edged interfaces*, J. Comput. Phys., 224 (2007), pp. 729–756.
- [41] Y. ZHOU, S. ZHAO, M. FEIG, AND G. WEI, *High order matched interface and boundary (MIB) schemes for elliptic equations with discontinuous coefficients and singular sources*, J. Comput. Phys., 231 (2006), pp. 1–30.